

Temperature

0th law of thermodynamics

There exists a useful quantity, temperature, which is a measure of the equilibrium state.

When two systems are in equilibrium ~~with~~, and there is no net energy flow between them, they have the same temperature.

Define temperature

$$\langle E_k \rangle \equiv \frac{3}{2} kT \quad \text{for an ideal gas}$$

$k = 8.617 \times 10^{-5} \text{ eV/K}$ - Boltzmann constant
and T is in Kelvin scale

for N_2 molecule

$$E_k = \frac{1}{2} m v^2$$

$$\left\langle \frac{m v^2}{2} \right\rangle = \frac{3}{2} kT$$

$$\sqrt{\langle v^2 \rangle} \sim 510 \text{ m/s}$$

Ideal gas law

$$\frac{F}{A} = \frac{\left(\frac{\Delta p}{\Delta t}\right)}{A}$$

from a single molecule $\Delta p = 2mv_x$

$$\frac{F}{A} = \frac{2mV_x}{A \Delta t}$$

only particles within $V_x \Delta t$ can hit within Δt

fraction of total volume that can hit area A

is
$$\frac{V_x \Delta t \cdot A}{V}$$

only $\frac{1}{2}$ are moving in the correct direction

$$H = \frac{V_x \Delta t A}{2V}$$

$$P_s = H \left(\frac{E}{A}\right) = \frac{mV_x^2}{V}$$

$$P = \frac{Nm \langle v_x^2 \rangle}{V}$$

$$\langle v^2 \rangle = \langle v_x^2 \rangle + \langle v_y^2 \rangle + \langle v_z^2 \rangle = 3 \langle v_x^2 \rangle$$

$$P = \frac{1}{3} \frac{Nm \langle v^2 \rangle}{V}$$

$$PV = \frac{2}{3} N \left\langle \frac{mv^2}{2} \right\rangle$$

$$\uparrow \equiv \frac{3}{2} kT$$

$$PV = NkT$$

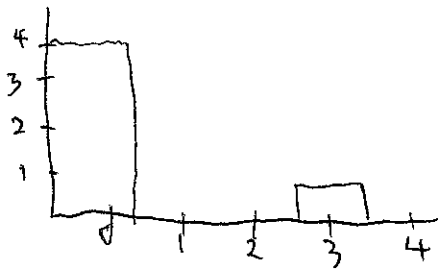
Maxwell - Boltzmann Distribution

If the total energy of a system is divided among a large number of particles, all possible divisions of energy are equally likely.

$$f_{MB} = C e^{-E/kT}$$

Example: total Energy = 3
5 particles

Combinations

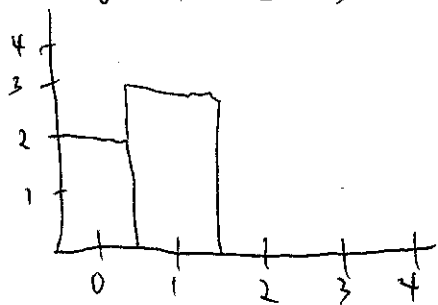


5

$\frac{\Delta N}{\Delta E}$

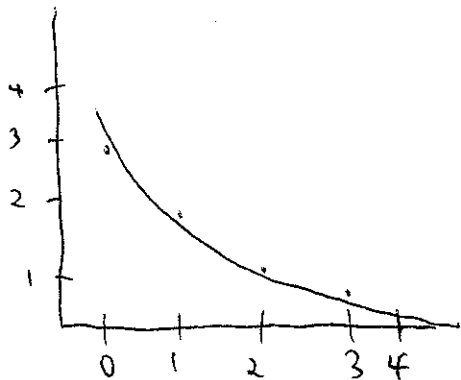


20



10

$\frac{dN}{dE}$



Another argument for exponential distribution

2 particles, total Energy = 10

$$E_1 + E_2 = 10$$

$P(E_1)P(E_2)$ - probability of particle 1 having $E = E_1$
and particle 2 having $E = E_2$

if all divisions of energy are equally likely

$$P(0)P(10) = P(1)P(9) = P(2)P(8)$$

$$P(E_1)P(E_2) = f(E_1 + E_2)$$

only $f = C e^{-E/KT}$

exponential can
do this

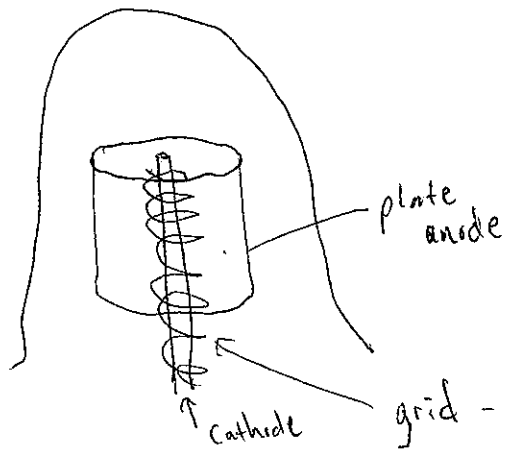
KT is denominator such that

$$\langle E \rangle = KT$$

correct behavior for ideal gas law is

restored by considering density of states.

Vacuum Tubes



$$J = C (kT)^2 e^{-\Phi/kT}$$

grid - can modulate current flow

Density of states

$$\rho(E) \equiv \frac{dn_s}{dE} \quad \text{- density of states}$$

$$n_s = \frac{N_s}{V} \quad \text{is number of states per volume with energy between } E \text{ and } E+dE$$

$$\frac{df}{dE} \propto \rho(E) e^{-E/KT} \quad \text{number of occupied states per volume with energy between } E \text{ and } E+dE$$

for an ideal gas

$$\rho(E) = \left(\frac{N}{V}\right) (2\pi) \left(\frac{1}{\pi kT}\right)^{3/2} \sqrt{E}$$

maxwell speed distribution

$$\frac{d^3f}{dv_x dv_y dv_z} = \left(\frac{m}{2\pi kT}\right)^{3/2} e^{-mv^2/2kT}$$

$$v = \sqrt{v_x^2 + v_y^2 + v_z^2}$$

$$dv_x dv_y dv_z = (dv)(v d\theta)(v \sin\theta d\phi)$$

$$\frac{d^3f}{dv d\theta d\phi} = v^2 \sin\theta \frac{d^3f}{dv_x dv_y dv_z}$$

$$\frac{df}{dv} = \int_0^\pi d\theta \sin\theta \int_0^{2\pi} d\phi v^2 \frac{d^3f}{dv_x dv_y dv_z}$$

$$\frac{df}{dv} = 4\pi \left(\frac{m}{2\pi kT}\right)^{3/2} v^2 e^{-mv^2/2kT}$$

Energy Distribution

$$E = \frac{1}{2} m v^2$$

$$dE = m v dv$$

$$\frac{df}{dE} = \frac{dv}{dE} \cdot \frac{df}{dv} = \frac{4\pi}{m v} \left(\frac{m}{2\pi kT} \right)^{3/2} v^2 e^{-mv^2/2kT}$$

$$\frac{df}{dE} = 2\pi \left(\frac{1}{\pi kT} \right)^{3/2} \sqrt{E} e^{-E/kT}$$

$$\frac{dn}{dE} = \frac{N}{V} \frac{df}{dE}$$

$$\frac{dn}{dE} = \rho(E) f_{MB}$$

$$\rho(E) = \left(\frac{N}{V} \right) (2\pi) \left(\frac{1}{\pi kT} \right)^{3/2} \sqrt{E}$$